

particle acceleration on Earth, ca. 1937

Particle acceleration in magnetized, astrophysical turbulent plasmas

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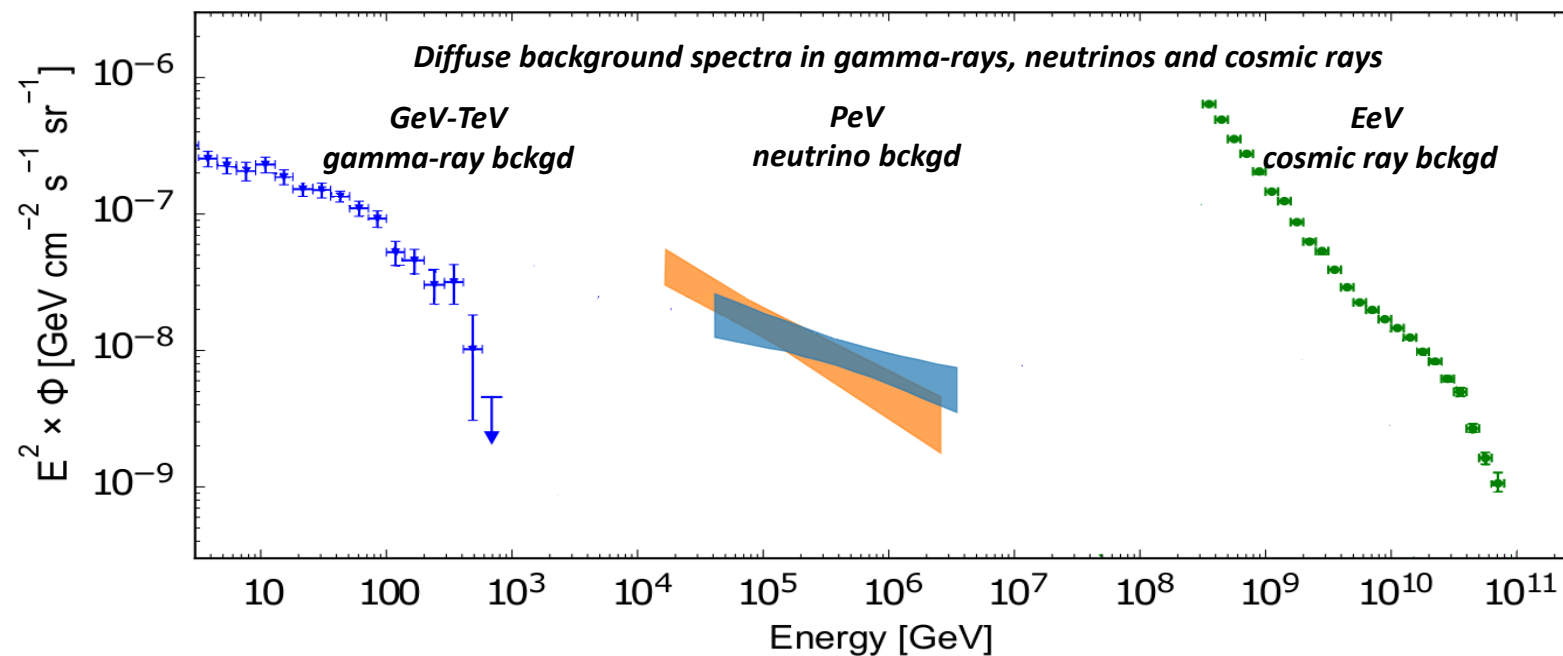
(CEA, France)

+ G. Pelletier (IPAG), A. Bykov (Ioffe), M. Malkov (UCSD),
L. Comisso & L. Sironi (Columbia)

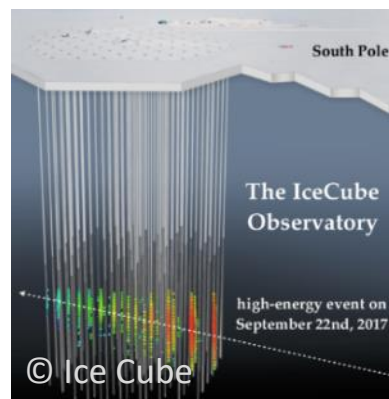
Physics of particle acceleration: a cornerstone of high-energy multi-messenger astrophysics

→ $\nu - \gamma - \text{CR}$ connection: acceleration of ions → cosmic rays, photons and neutrinos [e.g. Waxman + Bahcall 97, 98]

→ *what are the accelerating machine(s) and the acceleration process(es) at work ?!*



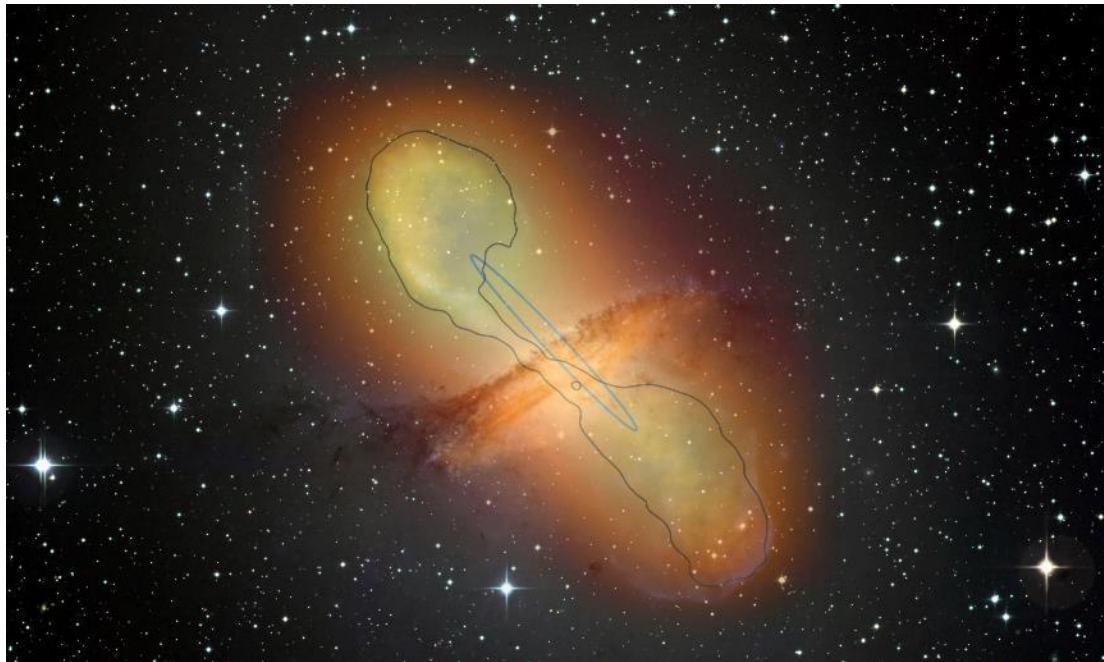
© Ice Cube, Ahlers + Halzen 17



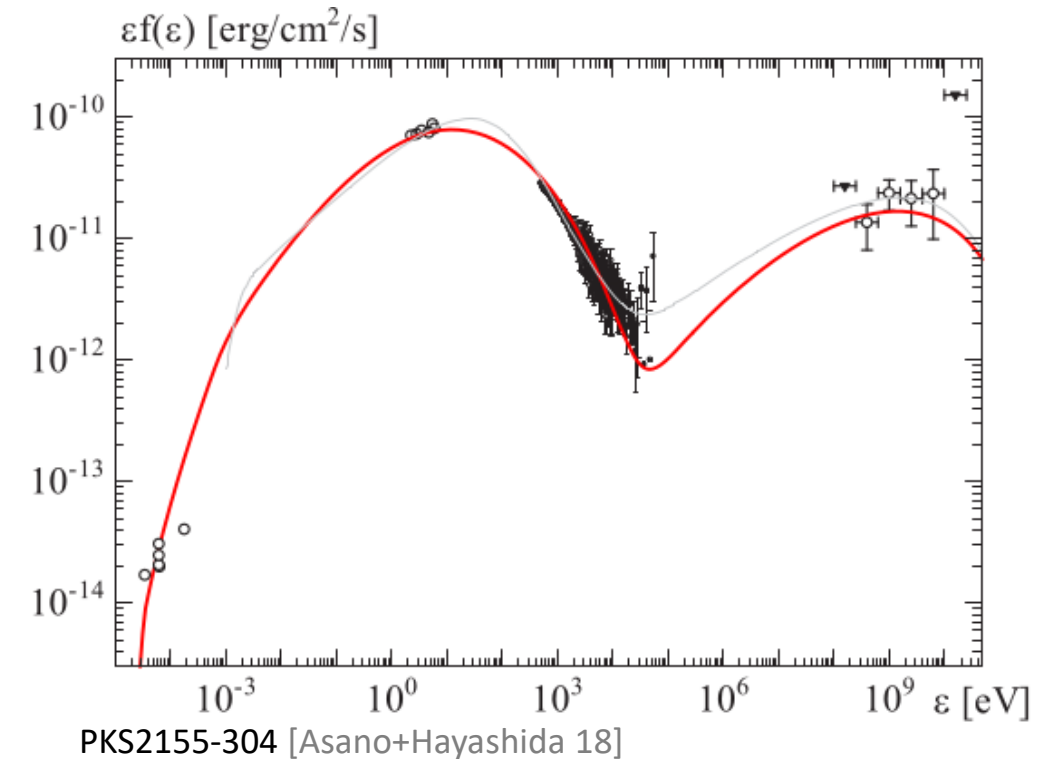
Stochastic Fermi acceleration on all scales: from large-scale jets to the blazar zone

→ in large-scale jets: continuous acceleration in turbulent flows invoked to explain the non-thermal emission seen on scales exceeding the cooling length scale of high-energy electrons... Refs.: e.g. Liu+17, Rieger 19, Webb+18,20

→ in blazars (radio-galaxies with jet head-on to observer): characteristic double-hump spectrum (synchrotron – inverse Compton) from non-thermal electrons... acceleration physics: reconnection, turbulence, shocks?

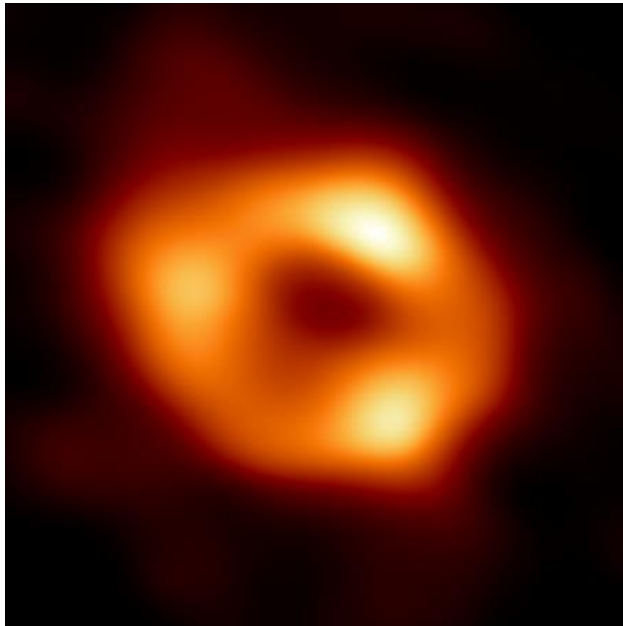


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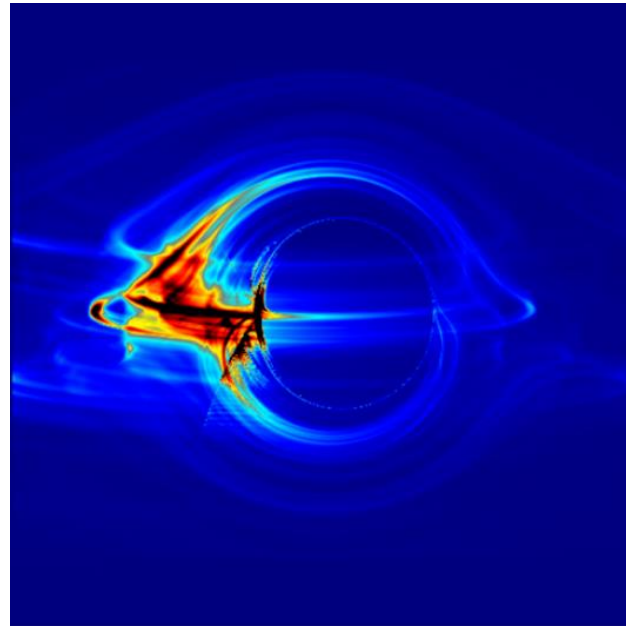


Stochastic Fermi acceleration on all scales: down to black hole surroundings

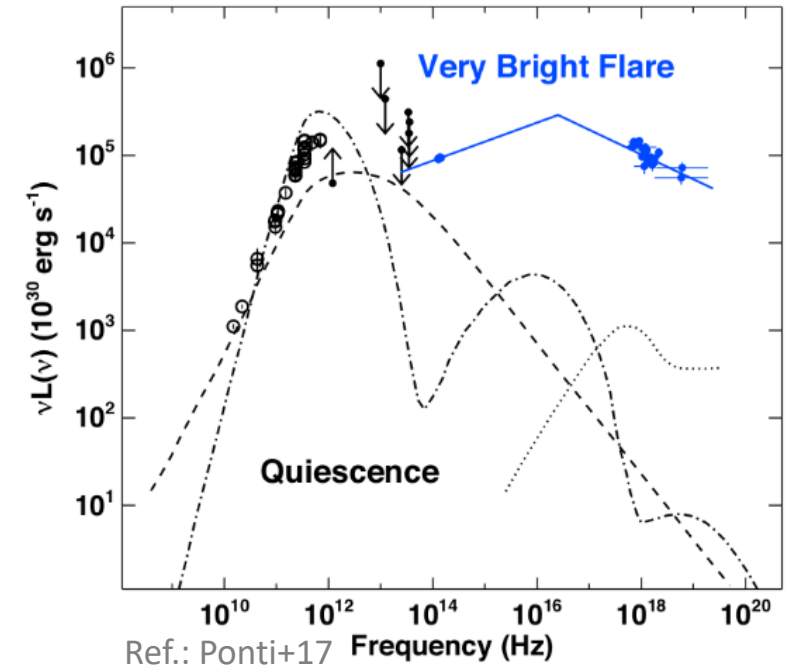
- Importance of microphysics: radiation of black hole flows shaped by the physics of dissipation in collisionless turbulence...
- Flares seen in NIR and X around SgrA*: suggest powerlaw extension with slope $\sim -3 \dots -2$, + synchrotron cooling ...
 - ⇒ key scenarios: reconnection (at large magnetization), or turbulence (if large fluctuations)?



Ref.: Event Horizon Telescope 22



Ref.: Petersen+Gammie20

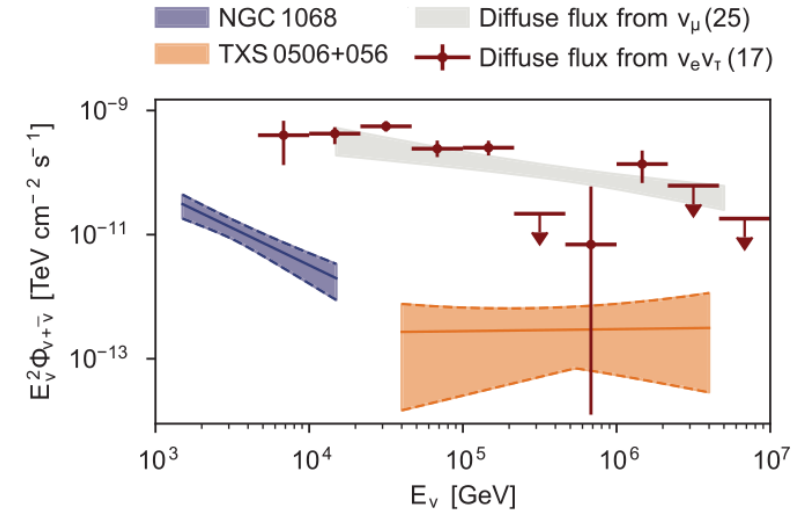
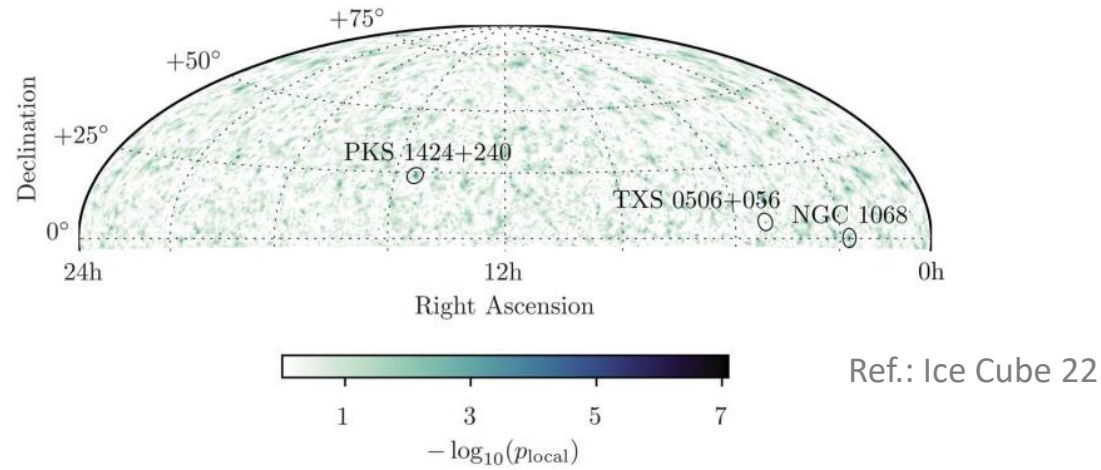


Ref.: Ponti+17

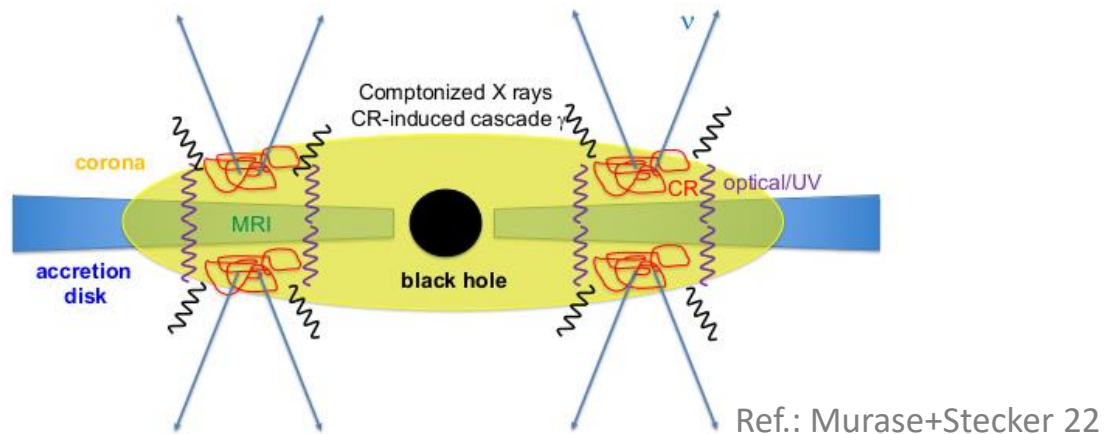
Note: EHT image well reconstructed by GRMHD simulations... which however use recipes to describe particle heating + acceleration!

Stochastic Fermi acceleration as an origin for the high-energy neutrinos from NGC1068

→ Ice Cube 22: a clear excess of high-energy (1-10 TeV) neutrinos from nearby AGN NGC7068



→ interpretation: e.g. [Murase+], particle acceleration in corona of accretion disk + neutrino production in $p - \gamma$ process



→ particle acceleration: turbulence, shocks?

Numerical studies of particle acceleration in collisionless magnetized turbulence

→ a non-linear, multi-scale problem:

... e.g. in turbulence: a fully nonlinear interplay between particles and e.m. fields...

⇒ HPC numerical simulations using « particle-in-cell » (PIC) method

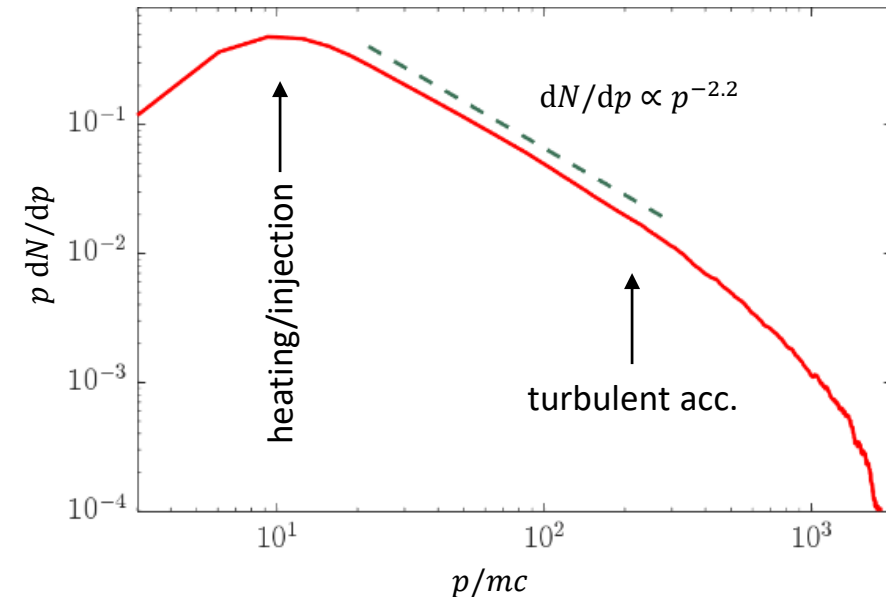
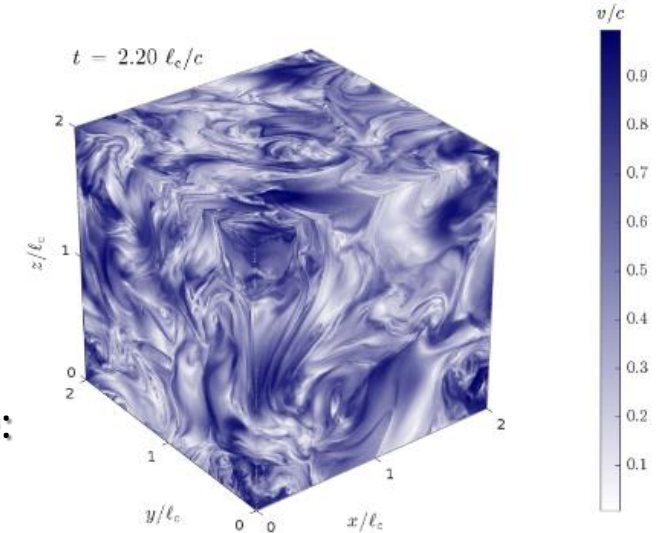
→ numerical experiments of particle acceleration in magnetized, collisionless turbulence:

... recent breakthrough results in (trans- and fully-) relativistic turbulence¹

... meaning: magnetization parameter $\sigma \equiv \frac{B^2/4\pi}{n m c^2} = (u_A/c)^2 \sim 1$

main findings:

1. two-stage particle acceleration:
first, injection through reconnection up to $p \sim \sigma m c$...
then, particle acceleration in “ideal” fields
2. unexpected² emergence of powerlaws
with $dN/dp \propto p^{-s}$ and $s \sim 2 \dots 4$



Refs: 1. Zhdankin+17,18,20,... Wong+ 19, Comisso+Sironi 18, 19, Nätilä + Beloborodov 20, ... Groselj+23 (+ many MHD sims)
2. discussion in M.L. + Malkov 20

Challenges in modeling particle acceleration from first-principles

→ a challenge of scales:

... microscopic acceleration scales:

... gyroradius: $r_g \sim 3 \times 10^6 \text{ cm } E_{\text{GeV}} B_G^{-1}$

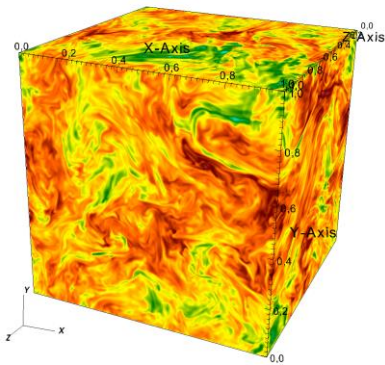
... macroscopic source scales:

... e.g. blazar zone $R \sim 10^{16} \text{ cm}$

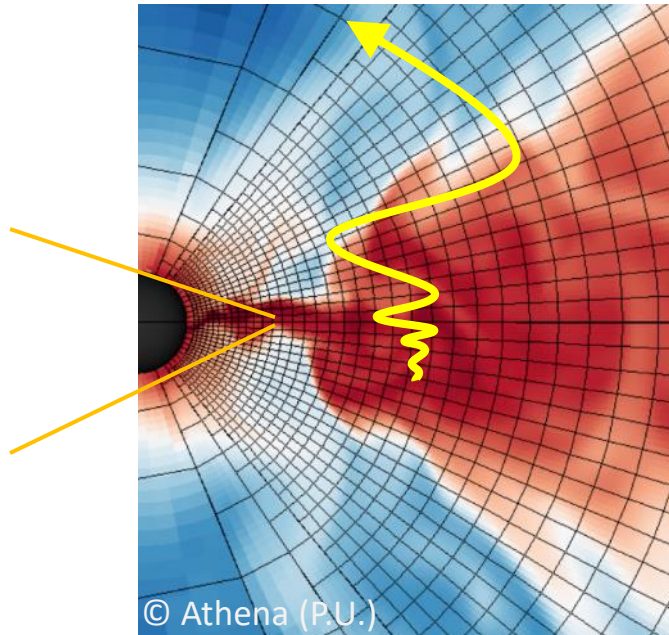
⇒ a strong limitation for applicability of PIC simulations: $10^3 \sim 2$ orders of magnitude in dynamic range..

... in practice: phenomenology (macro → micro) vs theory (micro → macro)

→ a need for microscopic recipes to model particle acceleration in complex, random velocity flows:



© C. Demidem, rel. MHD turb.



© Athena (P.U.)

→ multi-stage acceleration:

... scattering m.f.p. increases with energy ⇒ particle probes different velocity flows as energy increases

... from non-ideal/reconnection → turbulence → sheared velocity flows

Particle acceleration in magnetized, astrophysical turbulent plasmas

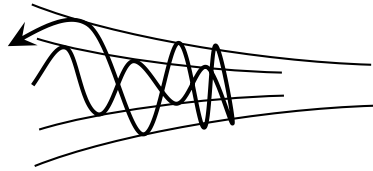
Outline:

1. General motivations and context
2. A colourized picture for stochastic Fermi acceleration:
 - contribution of non-resonant acceleration supported by numerical simulations
 - physics of acceleration shaped by the intermittency of turbulence
3. Discussion + remarks toward phenomenology



The Fermi picture for particle acceleration (1949, 1954)

→ assumption: perfectly conducting magnetized plasma composed of moving scattering centers...
 particle acceleration on motional electric fields $\mathbf{E} = -\mathbf{v}_E \times \mathbf{B}/c$



Fermi type A reflection of a cosmic-ray particle

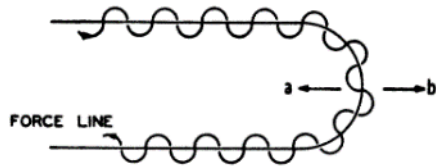
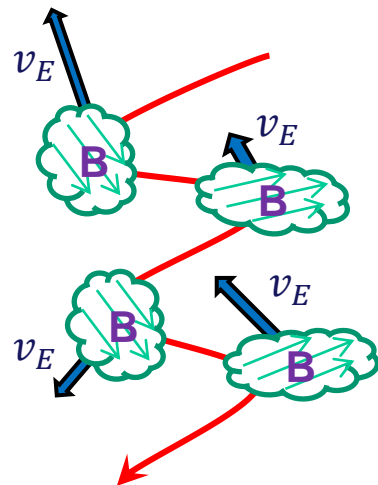


FIG. 1. Type B reflection of a cosmic-ray particle.



→ sequence of discrete interactions with point-like scattering centers... in each scattering center rest frame: elastic collision
 (ideal MHD $\Rightarrow \mathbf{E} = 0$ in rest frame)

→ kinematics: two-body collision, isotropic + elastic scattering in scattering center rest frame
 $\Rightarrow \Delta p > 0$ for head-on, $\Delta p < 0$ tail-on

→ stochastic acceleration (diffusion in momentum space)...
 e.g. Fokker-Planck equation:

$$\frac{\partial}{\partial t} f(p, t) = \frac{1}{p^2} \frac{\partial}{\partial p} \left[p^2 D_{pp} \frac{\partial}{\partial p} f(p, t) \right]$$

momentum diffusion coefficient: $D_{pp} \sim \frac{v_E^2}{c^2} \frac{p^2}{t_{\text{int}}}$

→ an issue: implementing stochastic acceleration in turbulence?

Generalized Fermi acceleration: implementation in a large-scale, random velocity flow

→ what matters is the shear of the velocity flow $\partial_\alpha u_E^\beta$:

ideal MHD conditions: \mathbf{E} vanishes in (comoving) frame moving at $\mathbf{u}_E \propto \mathbf{E} \times \mathbf{B}$

⇒ no acceleration in absence of shear...

... in original Fermi scenario:

shear \leftrightarrow difference in velocity of scattering centers

... in turbulent flow:

$\partial_\alpha u_E^\beta \supset$ compression, shear, vorticity...
with contributions from all scales of cascade...

→ follow the particle momentum in the (non-inertial) frame where $E = 0$ ¹:

in that frame, no electric field...

⇒ momentum variation \propto non-inertial forces characterized by velocity shear of \mathbf{u}_E

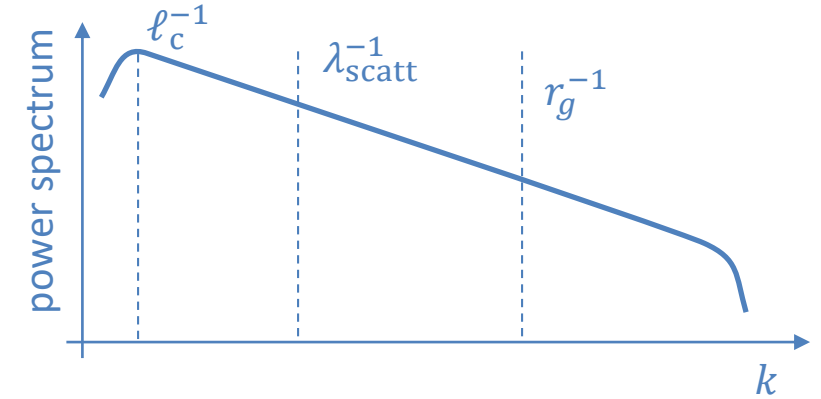
... \sim instantaneous Lorentz transform to non-inertial frame where interaction with e.m. field is elastic

(+ mandatory in relativistic settings)

Generalized Fermi acceleration: interaction with large scale modes

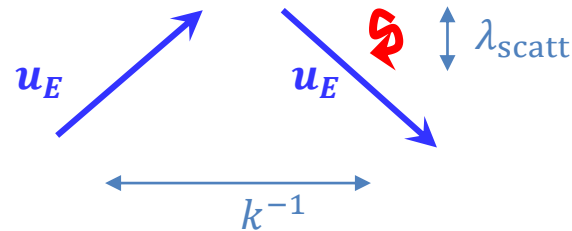
→ what (also) matters is how a particle experiences different scales:

... e.g.: for $r_g/\ell_c \rightarrow 0$ adiabatic limit (MHD)
 for $r_g/\ell_c \rightarrow \infty$ decoupling from turbulence



... for large scale modes: k (mode wavenumber) $\ll \lambda_{scatt}^{-1}$ (scattering m.f.p.) $\ll r_g^{-1}$ (gyroradius)

⇒ momentum diffusion coefficient¹:
$$D_{pp} \sim \underbrace{\frac{u_E^2}{c^2}}_{\text{Fermi scaling}} \underbrace{\frac{p^2}{\lambda_{scatt}/c} \left(\frac{\lambda_{scatt}}{k^{-1}} \right)^2}_{\ll 1}$$



... u_E gradient weak on scale λ_{scatt}
 ⇒ inefficient acceleration in comoving frame...

⇒ “shielding” from large scale modes $k \ll \lambda_{scatt}^{-1}$

Generalized Fermi acceleration: interaction with intermediate scale modes

→ dominant contribution: intermediate-scale modes with $\lambda_{\text{scatt}}^{-1} < k \leq \text{gyroradius } r_g^{-1}$

... transport \sim gyration around local magnetic field lines, i.e. coarse-grained on scale r_g

→ model²:

$$\frac{d\gamma'}{d\tau} = -\gamma' u'_{\parallel} \mathbf{a}_E \cdot \mathbf{b} - u'_{\parallel}{}^2 \Theta_{\parallel} - \frac{1}{2} u'_{\perp}{}^2 \Theta_{\perp}$$

energy change
in local comoving
frame

$$u'_{\parallel} = \mathbf{p}' \cdot \mathbf{b} / mc$$

$$u'_{\perp} = [u'^2 - u'_{\parallel}{}^2]^{1/2}$$

effective gravity
along field line

$$\mathbf{a}_E = u_E^{\alpha} \partial_{\alpha} \mathbf{u}_E$$

velocity shear
along field line

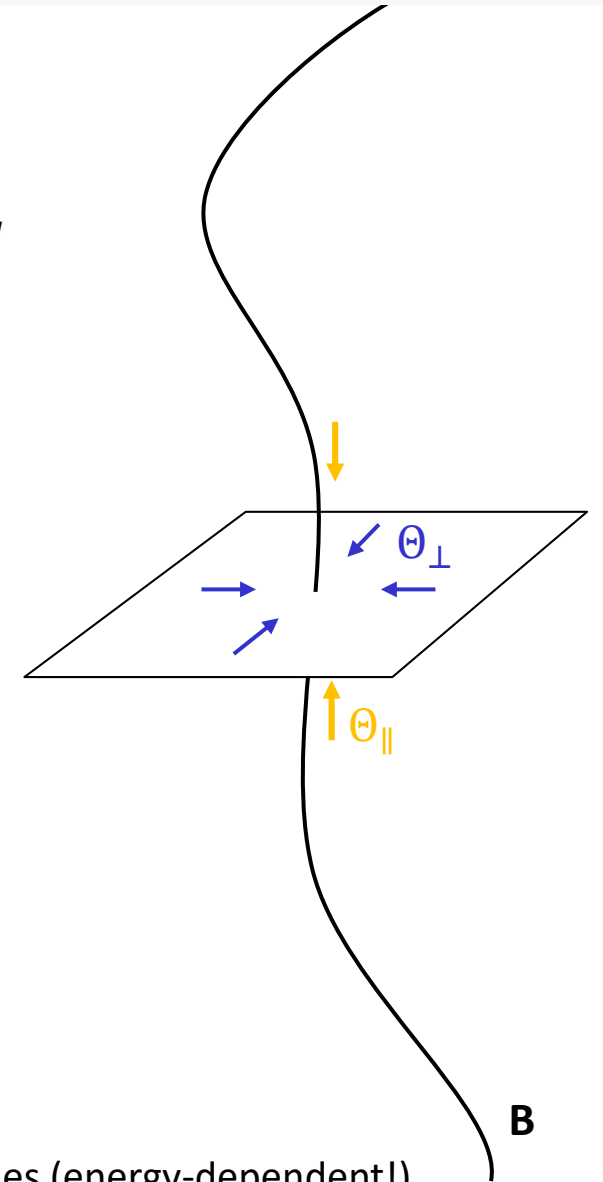
$$\Theta_{\parallel} = b^{\alpha} b^{\beta} \partial_{\alpha} u_{E\beta}$$

[Fermi type-B]
[field line curvature]

compression transverse
to field line

$$\Theta_{\perp} = (\eta^{\alpha\beta} - b^{\alpha} b^{\beta}) \partial_{\alpha} u_{E\beta}$$

[Fermi type-A]
[magnetic mirrors]



→ remarks:

... terms $\mathbf{a}_E \cdot \mathbf{B}$, Θ_{\parallel} and Θ_{\perp} are random forces: \Rightarrow random walk in momentum space
 \Rightarrow provides the required generalization of Fermi model to turbulent modes...

... average over gyro-orbit: \sim drift-kinetic theory in magnetic field coarse-grained on r_g scales (energy-dependent!)

Non-resonant Fermi-type acceleration in velocity gradients: distinctive features

→ acceleration scales with gradient of magnetic energy density

... unlike quasi-linear theory: \propto magnetic energy density

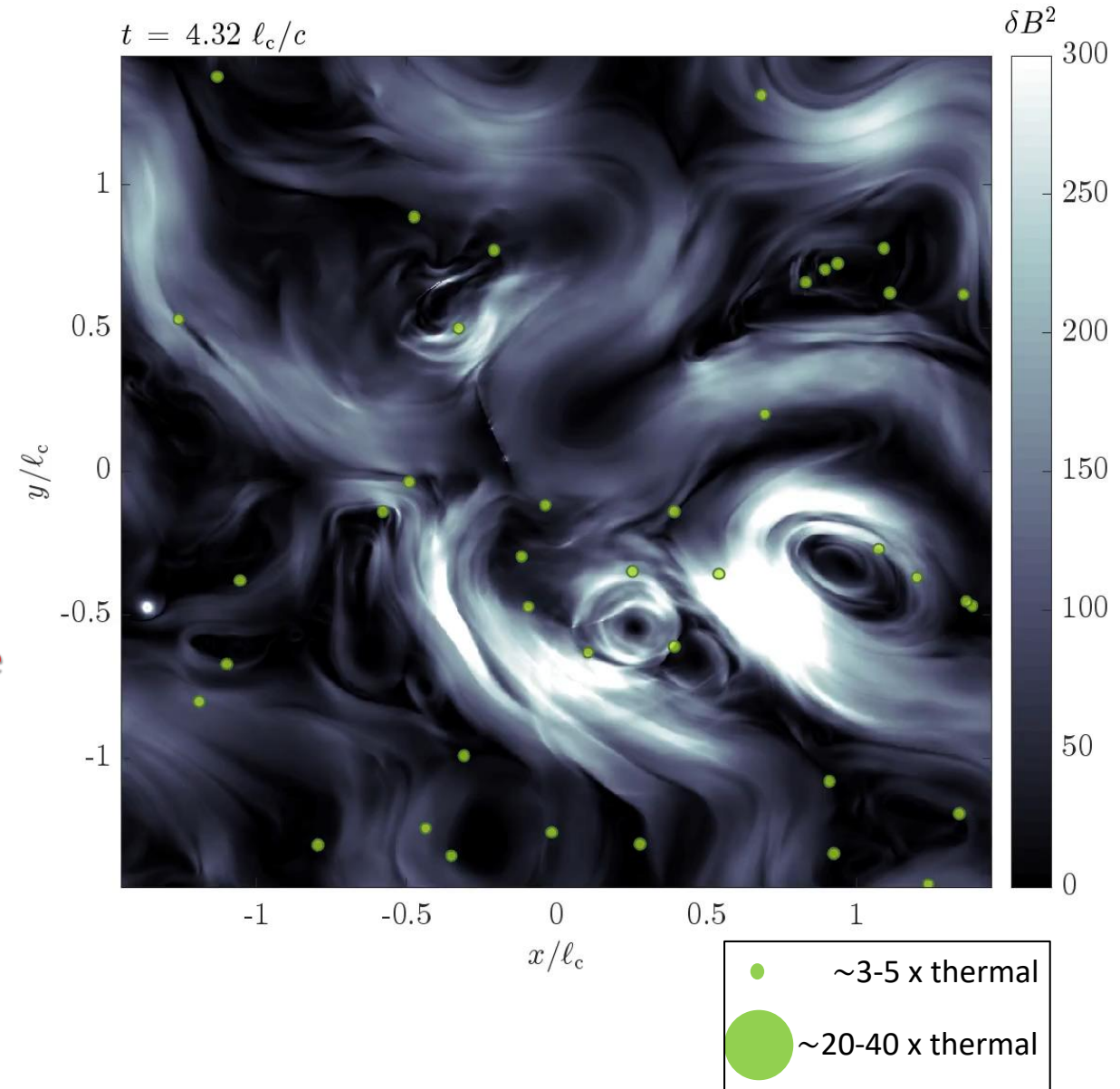
→ in each site, particle gains or loses energy regularly, according to sign of gradient

... unlike Fermi: head-on vs tail-on

→ acceleration sites occupy only a small filling fraction of the total volume

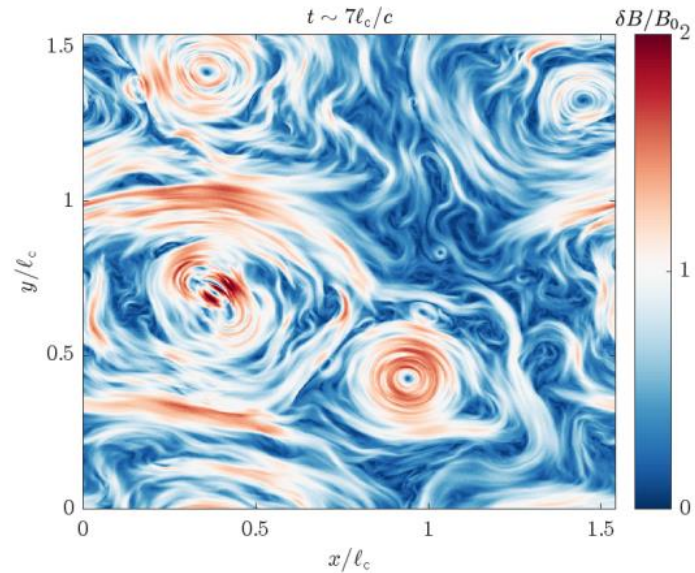
⇒ direct connection to intermittency

... unlike quasi-linear theory: homogeneous statistics

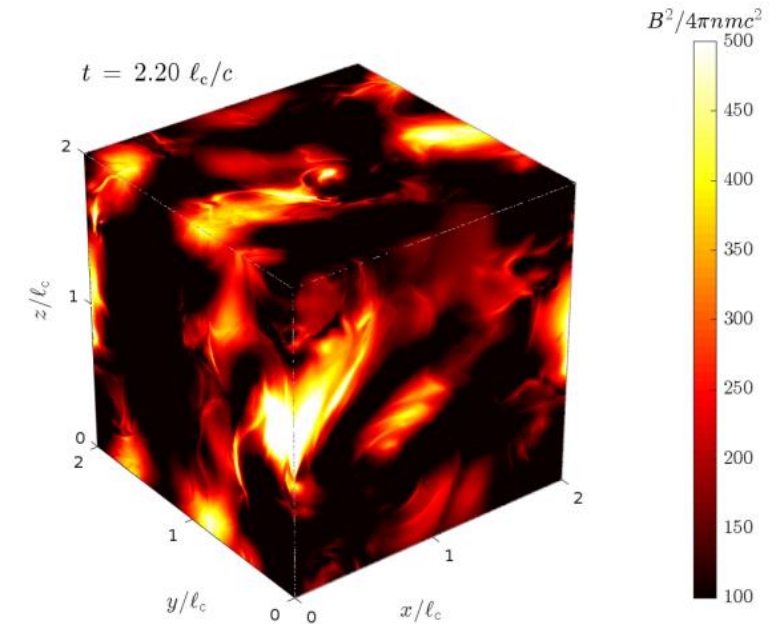


Non-resonant Fermi-type acceleration: comparison to numerical experiments

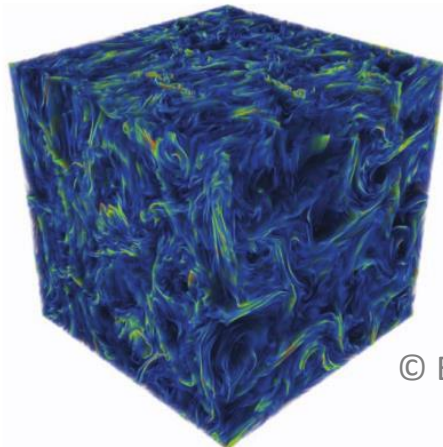
2D PIC simulation¹: forced and decaying, $10\,000^2$, e^-e^+ , $\delta B/B \sim 3$, $\sigma \sim 1$



3D PIC simulation¹: forced, $1\,080^3$, e^-e^+ , $\delta B/B \sim 3$, $\sigma \sim 1$



3D MHD simulation²: forced, $1\,024^3 \times 1\,024$, $\langle B \rangle = 0$, $v_A / c = 0.4$



© Eyink+13

+ synthetic turbulence³: sum of plane waves (Alfvén or fast magnetosonic)

Note: magnetization parameter $\sigma = \frac{\text{mag. energy dens.}}{\text{plasma energy dens.}}$

Non-resonant Fermi-type acceleration: comparison to numerical experiments

→ model:

$$\frac{d\gamma'}{d\tau} = -\gamma' u'_{\parallel} \mathbf{a}_E \cdot \mathbf{b} - u'_{\parallel}{}^2 \Theta_{\parallel} - \frac{1}{2} u'_{\perp}{}^2 \Theta_{\perp}$$

$$u'_{\parallel} = \mathbf{p}' \cdot \mathbf{b} / mc$$

$$u'_{\perp} = [u'^2 - u'_{\parallel}{}^2]^{1/2}$$

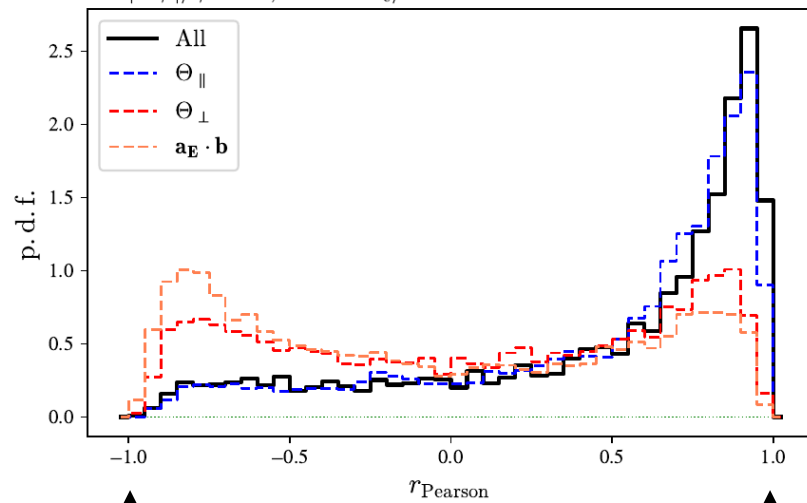
→ test¹:

for each particle history in a simulation, reconstruct $\gamma'(t)$ using above model and velocity gradients measured in the simulation at $\mathbf{x}, t...$

... then measure degree of correlation r_{Pearson} between the observed and reconstructed $\gamma'(t)$

PIC simulation: 3D, $1\,080^3$, e^-e^+ , $\delta B/B \sim 3$, $\sigma \sim 1$

$|\Delta\gamma'|/\gamma' > 2.0$, $\Delta t = 3.7 \ell_c/c$

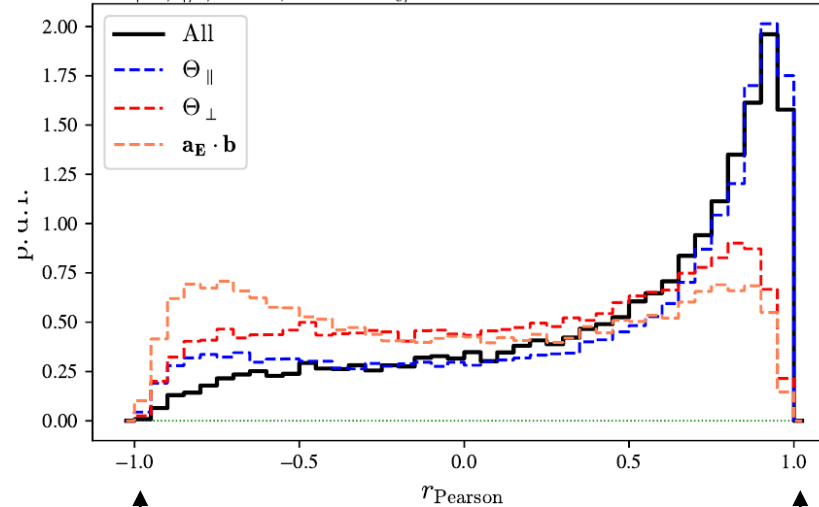


100% anti-correlation

100% correlation

PIC simulation: 2D, $10\,000^2$, e^-e^+ , $\delta B/B \sim 3$, $\sigma \sim 1$

$|\Delta\gamma'|/\gamma' > 1.0$, $\Delta t = 1.0 \ell_c/c$

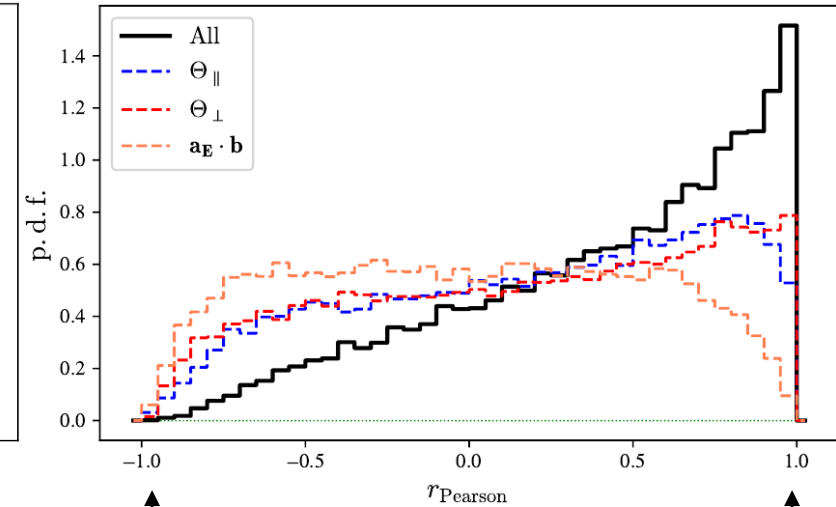


100% anti-correlation

100% correlation

driven incompress. MHD², 3D, 1024^3 , $v_A = 0.4c$

$|\Delta\gamma'|/\gamma' > 0.50$, $\Delta t = 1.70 \ell_c/c$



100% anti-correlation

100% correlation

⇒ **model captures the dominant contribution to particle energization**

+ note: in wave turbulence w/ resonant wave-particle interactions, no apparent correlation seen (as expected)

Powerlaw spectra: shaped by the intermittency of turbulence

→ statistics of the random force (\sim velocity gradient):

... velocity gradients become increasingly non-Gaussian (intermittent) at small scales (\leftrightarrow small gyroradii), taking large values in localized regions...

→ particle acceleration:

... some particles interact frequently with strong scattering centers, some not at all, even over long timescales ...

⇒ **anomalous transport¹ + powerlaws in momentum**

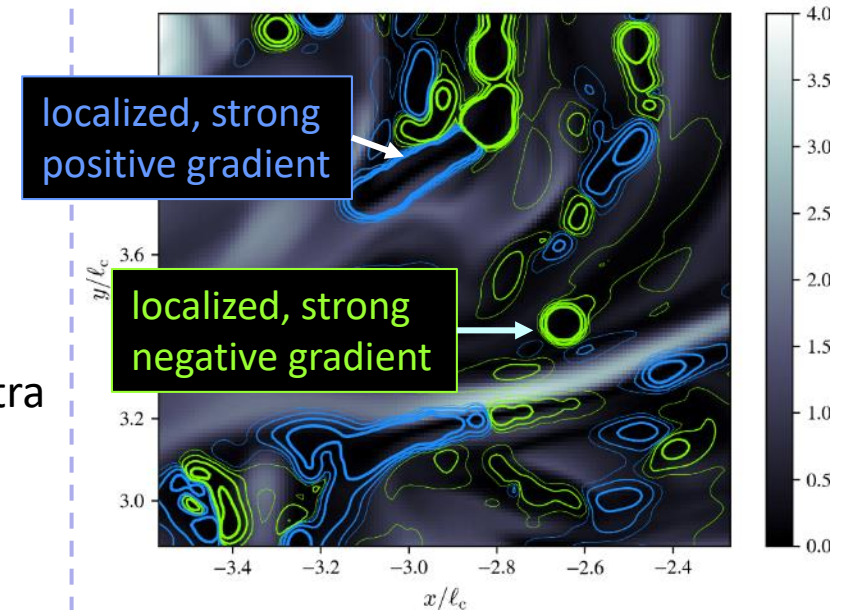
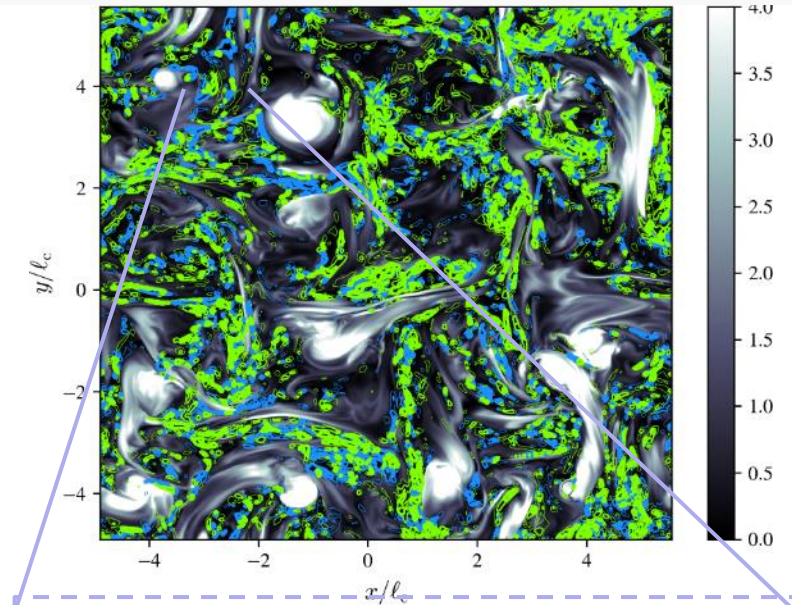
→ transport equation for distribution function²:

... failure of Fokker-Planck³: noise is non-Gaussian + non-white noise...

... derivation of a new transport equation:

pdf(momentum jump) \sim intermittency statistics

... transport equation produces powerlaws, accounts for particle spectra from time-dependent tracking in MHD simulation



3. see however Wong+19, Zhdankin+20

A transport equation for non-resonant particle acceleration in intermittent turbulence

→ note: random forces $(a_E \cdot b, \Theta_{\parallel}, \Theta_{\perp}) \sim \Gamma_l \neq$ Gaussian white noise
 \Rightarrow transport equation deviates from Fokker-Planck...

intermittency \sim origin of powerlaw

→ scheme: random force Γ_l (coarse-grained on scale $l \sim r_g$), p.d.f. $\text{Prob}(\Gamma_l)$

\Rightarrow momentum p jumps on timescale $\sim l/c$ by

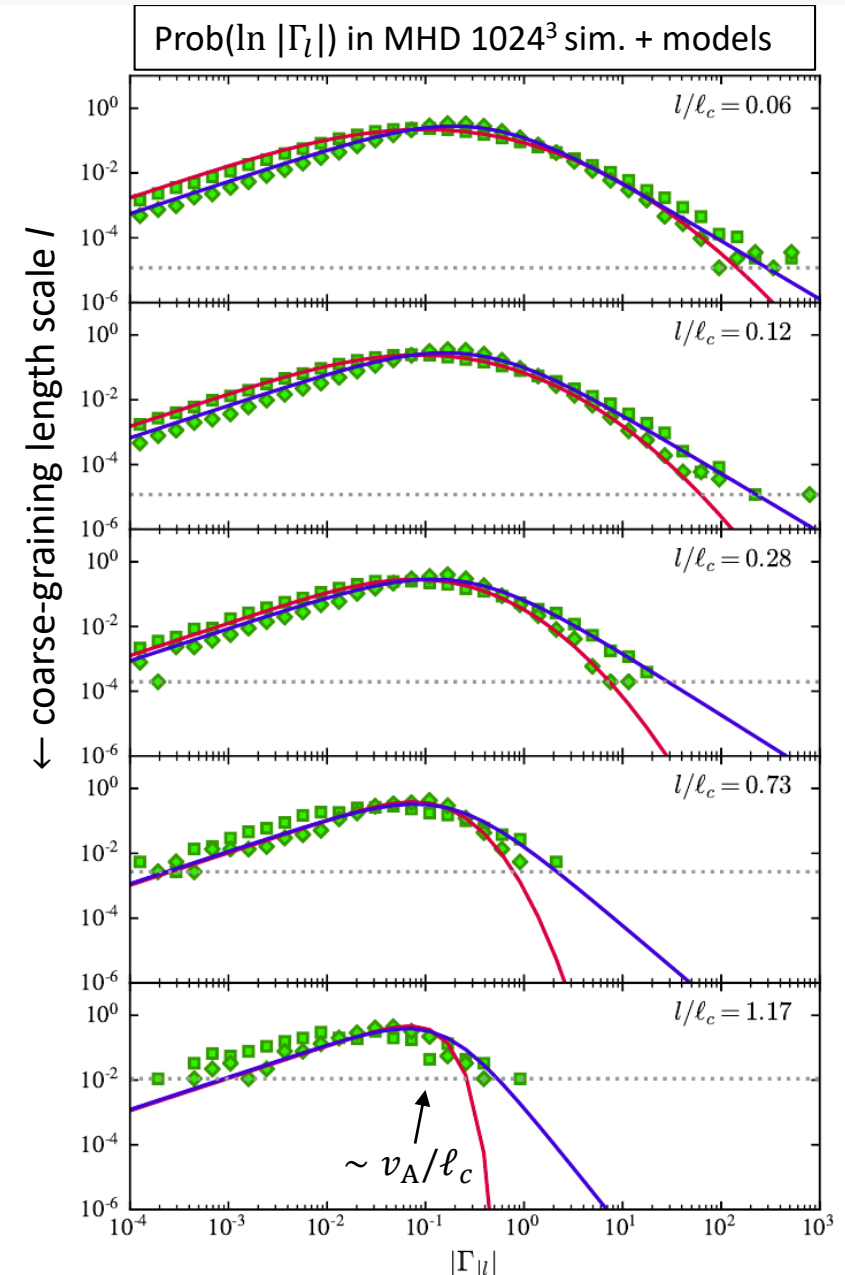
$$\Delta \ln p \sim \Gamma_l \Delta t \Rightarrow \text{Prob.}(\Delta \ln p) \sim \text{Prob.}(\Gamma_l)$$

→ kinetic equation¹:

$$\partial_t n_p = \int_0^{+\infty} dp' \left[\frac{\varphi(p|p')}{t_{p'}} n_{p'}(t) - \frac{\varphi(p'|p)}{t_p} n_p(t) \right]$$

$n_p = \frac{dN}{dp}$ $t_p \sim l/c$

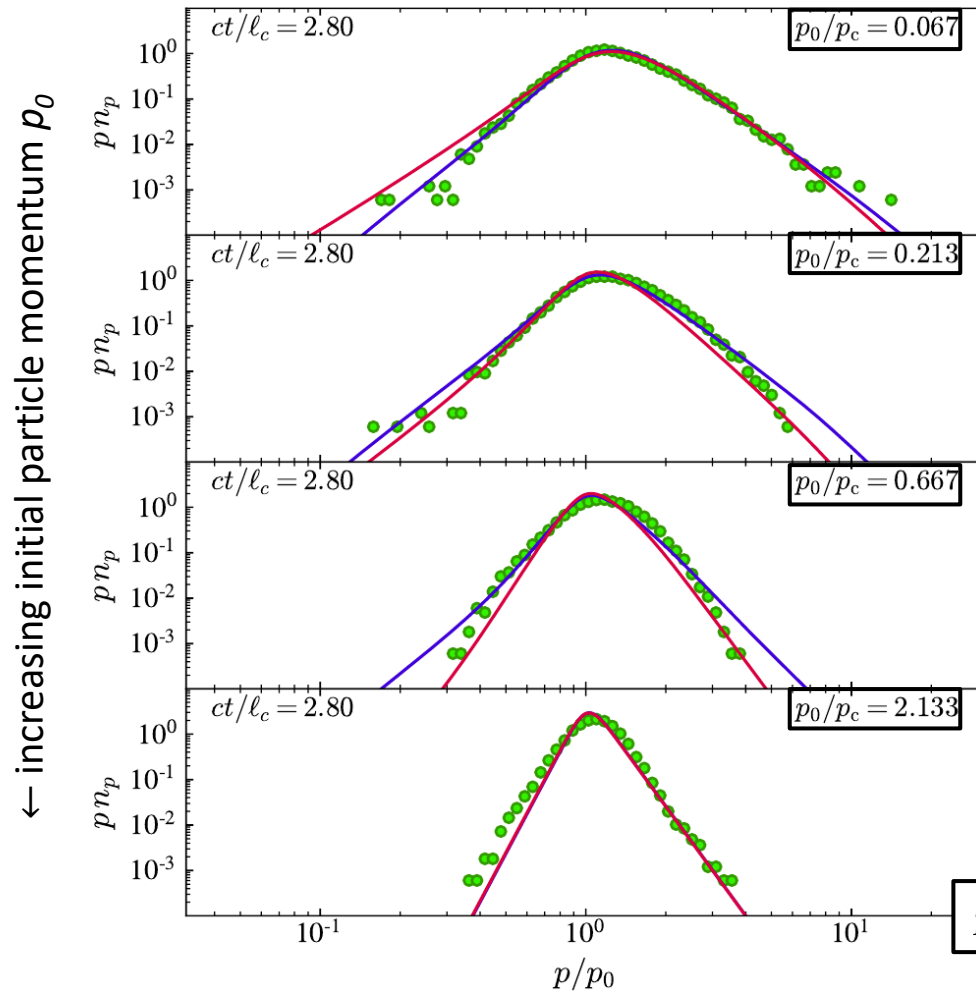
... hope: gain fundamental knowledge on $\text{Prob}(\Gamma_l)$ to model acceleration (e.g. intermittency studies²)



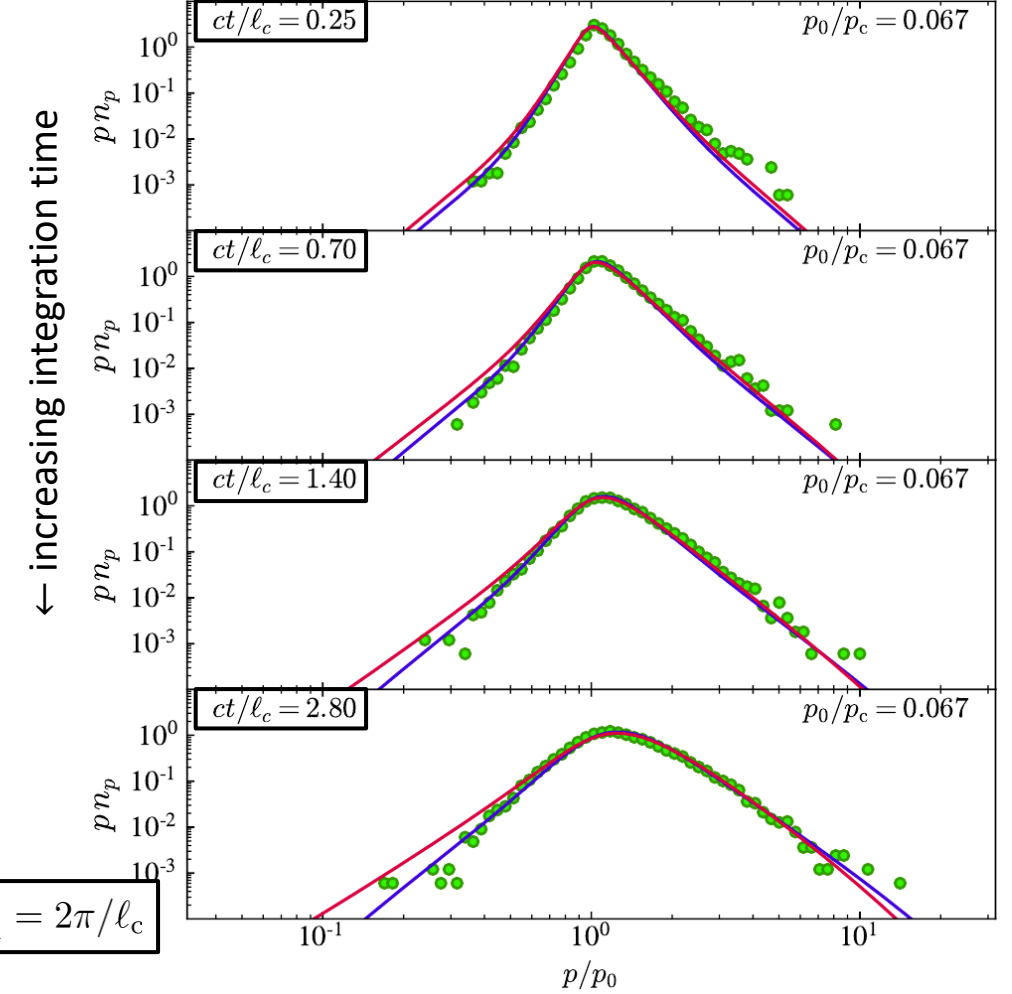
Accelerated particle spectra from phase space transport in intermittent turbulence

→ comparison to numerical data:

integrate kinetic equation and compare solution (Green function) to distribution measured in MHD 1024³ simulation by time-dependent particle tracking...



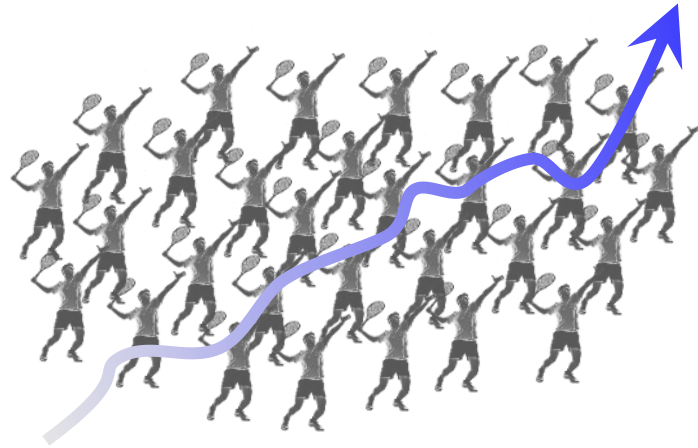
$$p_c : r_g(p_c) = k_{\min}^{-1} = 2\pi/\ell_c$$



⇒ transport equation can reproduce time- and energy- dependent Green functions... + capture powerlaw spectra
 ... in sub-relativistic regime: $dN/dp \propto p^{-4}$ and acceleration timescale $\sim \ell_c/v_A^2$ [see also Comisso + Sironi 22]

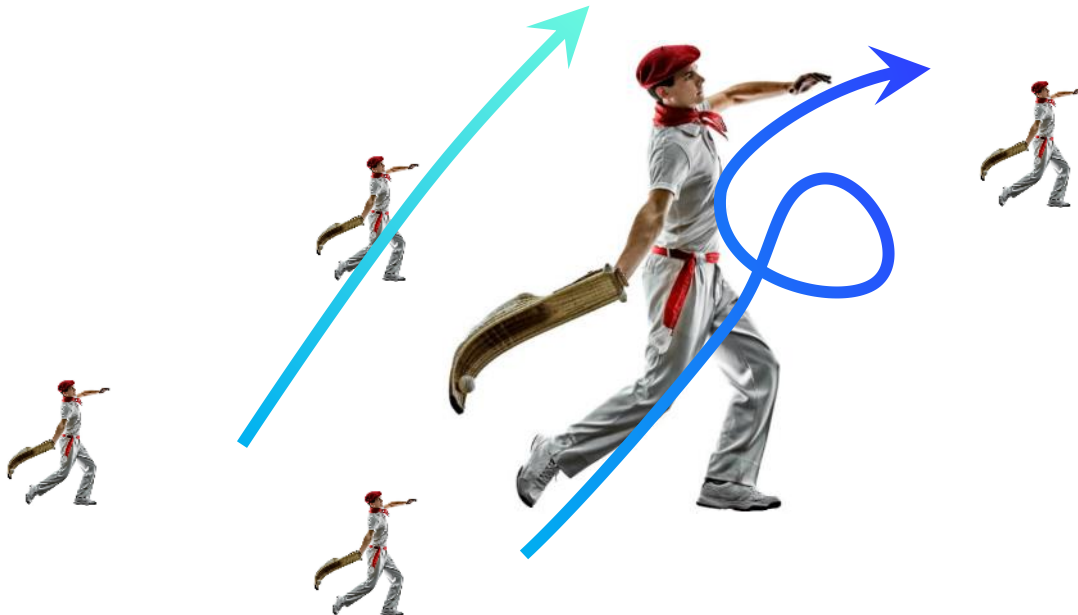
Colourizing the Fermi picture...

→ the original picture: stochastic acceleration as Brownian motion...



Brownian motion \leftrightarrow Fokker-Planck description,
characterized by one diffusion coefficient D_{pp} (+advection)

→ the colourized picture: stochastic interactions with *intermittent gradients*...



one diffusion coefficient D_{pp} does not describe spectra...
... particle acceleration dominated by intermittency...
... spectra exhibit powerlaw shapes...
... dominant acceleration: field line curvature...

Summary + discussion

→ Implementing (non-resonant) Fermi-type acceleration in a realistic turbulence setting:

... track particle history in frame in which $\mathbf{E}=\mathbf{0}$...

... particles are accelerated in regions of strong velocity gradients

→ Test on PIC + MHD simulations: the Fermi picture is well alive

... model captures bulk of energization in supra-thermal powerlaw region... at $\sigma \gtrsim 0.1$

→ Deriving a transport equation for Fermi acceleration:

... velocity gradients are non-Gaussian on small scales: intermittency rules...

... a multi-fractal model of gradient statistics, and a transport equation...

→ Some limitations:

... extrapolation to small spatial length scales ?

... role of turbulence anisotropy, particle trapping in structures?

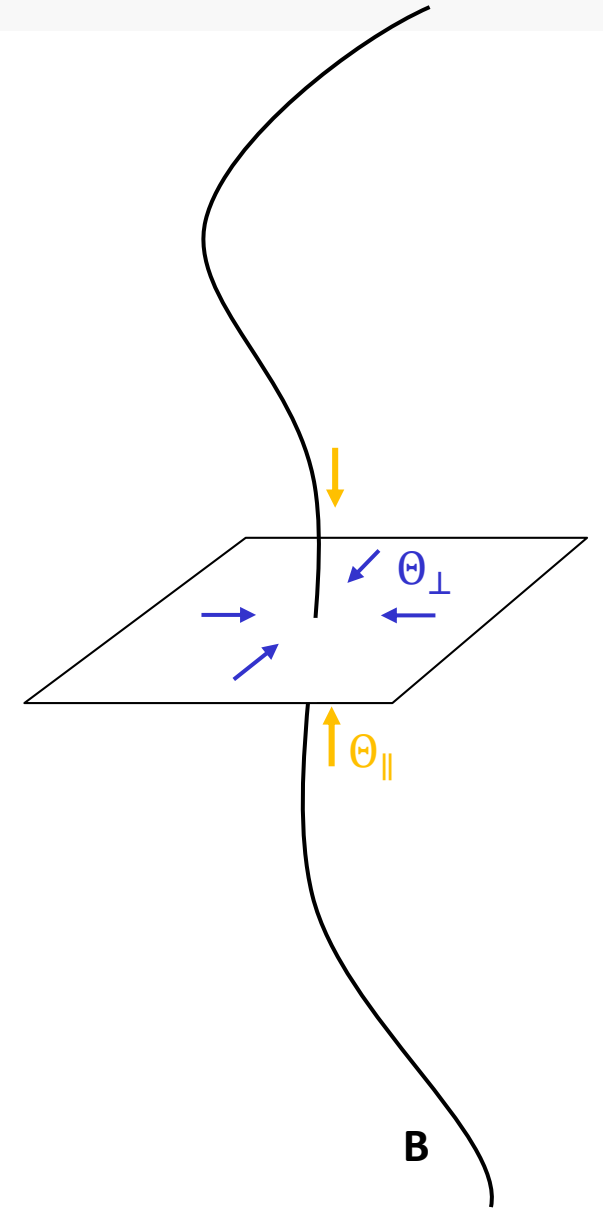
→ Some perspectives:

... better understanding the role and nature of intermittency wrt acceleration...

... consequences for phenomenology: flares etc...

... generalization to transport: e.g., role of intermittent magnetic mirrors...

... recipes for incorporating particle phase space transport in large-scale numerical simulations?



Statistics of turbulence intermittency

→ intermittent statistics: e.g. $\langle |\delta u_l^2| \rangle \propto l^{2h}$ but $\langle |\delta u_l|^n \rangle \not\propto l^{nh}$ ($n > 2$)

... structure functions: $S_n \equiv \langle |\delta u_l|^n \rangle = l^{\zeta_n}$

ζ_n in one-to-one correspondence with p.d.f. of δu_l

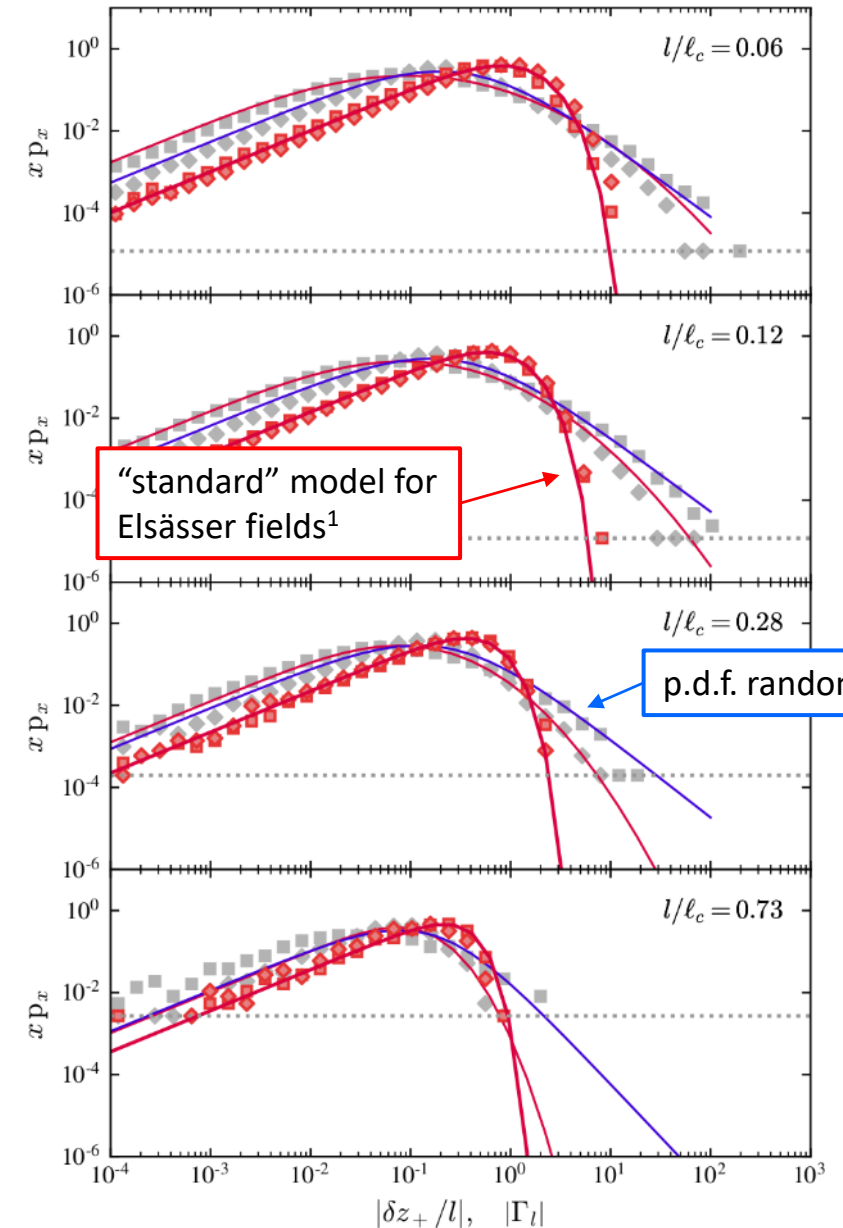
→ intermittent statistics of random forces:

... more extended, broader powerlaw tails than “standard” statistics for Elsässer fields $\delta z_{\pm} \equiv \delta v_l \pm \delta b_l$

... p.d.f. of random forces connected to field line curvature, which displays powerlaw behavior²

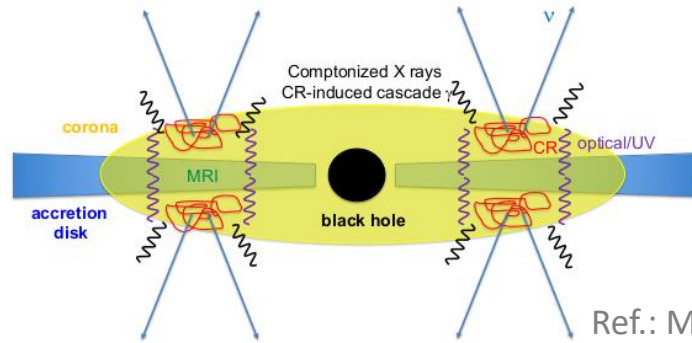
... a model for statistics of field line curvature: see ³

→ in different types of turbulence (e.g. compressive vs incompressible), different spectra⁴ because of different statistics?

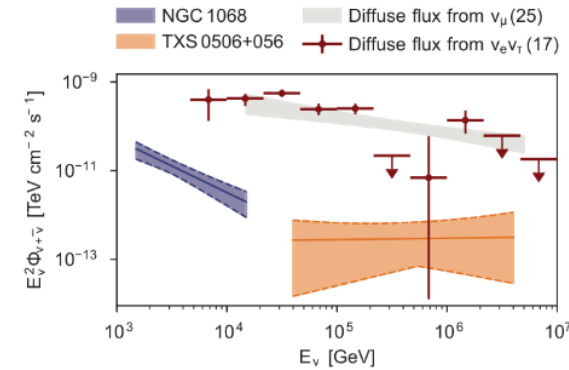


Stochastic Fermi acceleration as an origin for the high-energy neutrinos from NGC 1068

→ Ice Cube 22: a clear excess of high-energy (1-10 TeV) neutrinos from nearby AGN NGC 1068...
 ... a possible scenario: stochastic acceleration in turbulent corona + $p - \gamma$ neutrino production¹



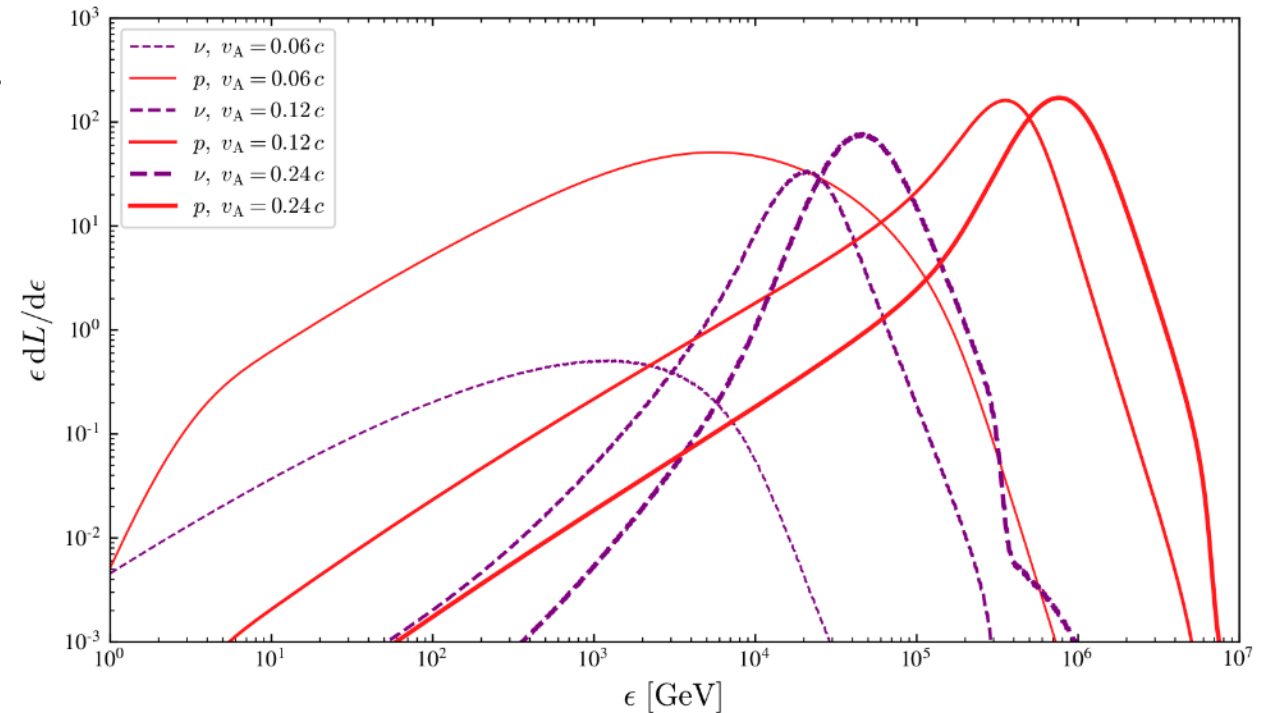
Ref.: Murase+Stecker 22



Ref.: Ice Cube 22

→ model: integrate spectra through transport eqn...
 ... including relevant energy losses¹

⇒ p acceleration to $>100\text{TeV}$ possible for
 turbulent Alfvén velocity $v_A \gtrsim 0.1c$
 ... +shear contribution?²



Refs.: 1. e.g. Murase 22 + refs.

2. M.L. + Rieger, in prep.

Does intermittency affect spatial transport?

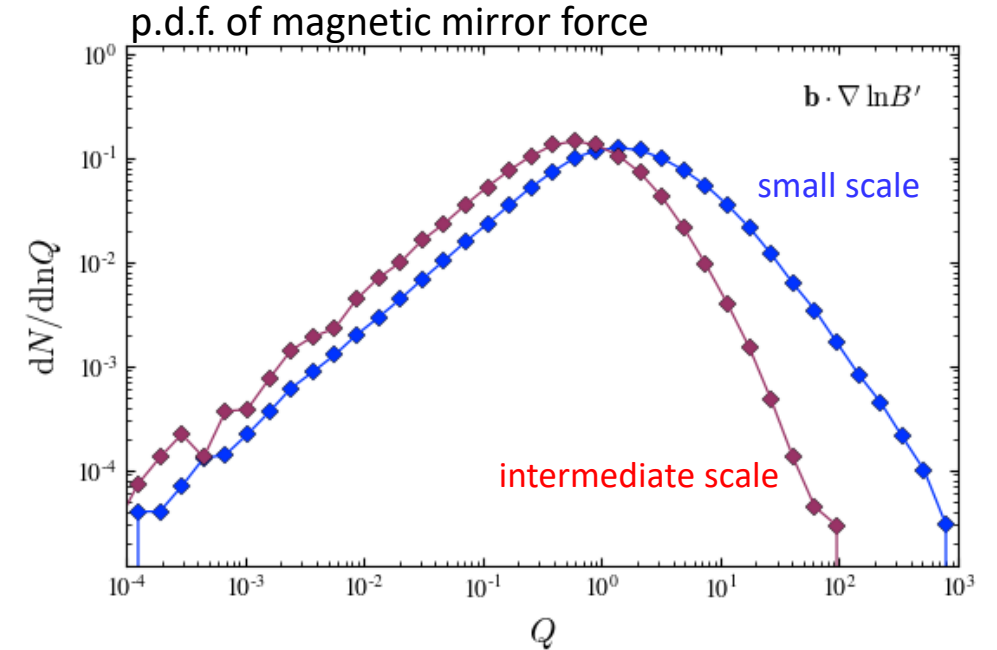
→ anomalous spatial transport by scattering on intermittent structures?

→ in absence of resonant wave-particle interactions: scattering ← magnetic mirrors¹

⇒ subject to intermittency...

⇒ expect spatial anomalous diffusion on scales $\sim \ell_c$:
... superdiffusion for some particles,
... trapping for others... seen in some simulations³

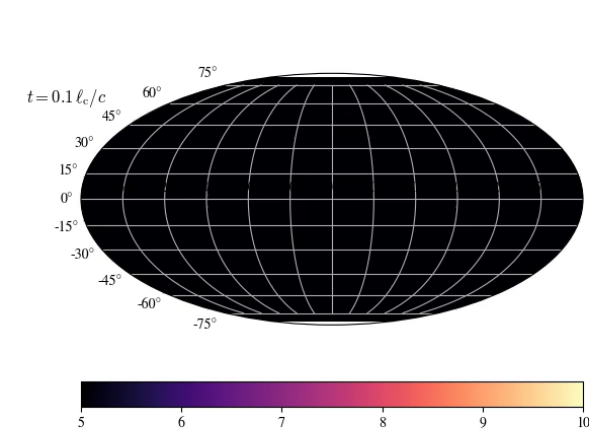
⇒ phenomenological consequences for e.g. pulsar halos,
cosmic-ray anisotropies at high energies etc.?



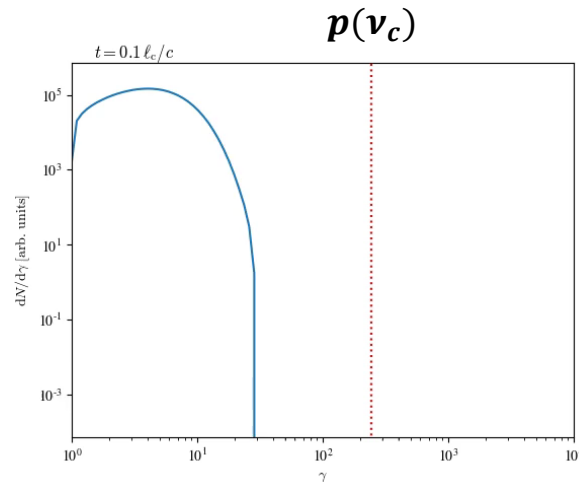
→ note: for low-energy cosmic rays, no intermittency effect because of very large travel time...

Consequence of intermittency for radiative signatures

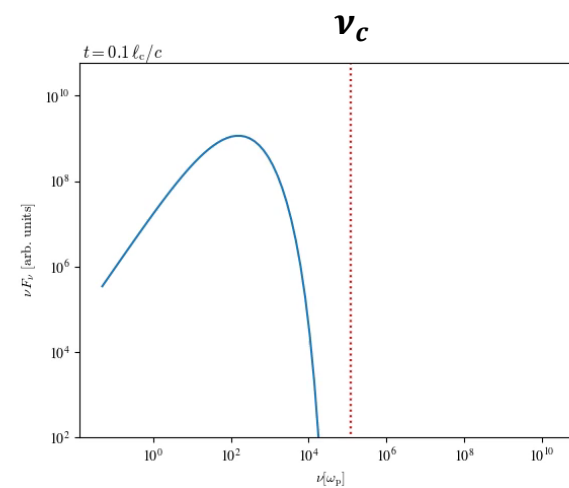
→ intermittency and high-energy flares: particle distribution highly anisotropic with spectral shape non-uniform in space close to the maximal energy¹...



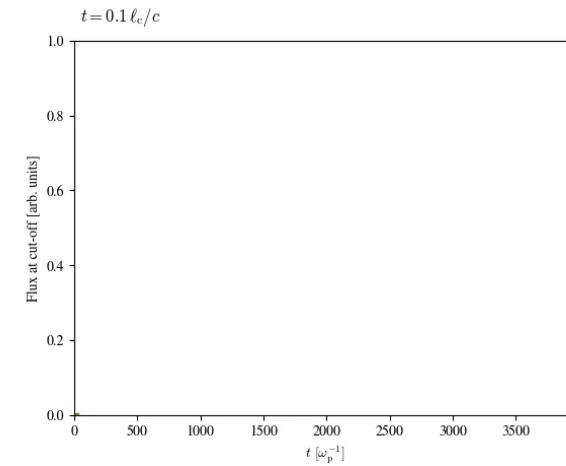
synchrotron map at cut-off ν_c



$p \frac{dN}{dp}$ vs p



νF_ν vs ν



light curve at cut-off ν_c

→ important phenomenological consequences for time-dependent flaring sources at high energies (e.g. blazars, GRBs etc.)

→ a potential realization of [Bykov+13] scenario for Crab flares: injection of high-energy pairs in turbulence?

→ note: in relativistic turbulence $\sigma \gtrsim 1$, acceleration timescale $t_{\text{acc}} \propto 1/u_E^2 \propto 1/\sigma$!

Summary and perspectives

→ Implementing (non-resonant) Fermi-type acceleration in a realistic turbulence setting:

... track particle history in frame in which $\mathbf{E}=\mathbf{0}$...

... particles are accelerated in regions of strong velocity gradients

→ Test on PIC + MHD simulations: the Fermi picture is well alive

... model captures bulk of energization in supra-thermal powerlaw region...

→ Deriving a transport equation for Fermi acceleration:

... non-Gaussian velocity gradients on small scales: intermittency rules...

... a multi-fractal model of gradient statistics, and a transport equation...

→ Some limitations:

... extrapolation to small spatial length scales ?

... role of turbulence anisotropy, particle trapping in structures?

→ Some perspectives:

... role and nature of intermittency wrt acceleration...

... consequences for phenomenology: flares etc...

... generalization to transport: e.g., role of intermittent magnetic mirrors...

... recipes for particle phase space transport in numerical simulations?

HEPRO VIII

High Energy Phenomena in Relativistic Outflows VIII

Paris, 23-26 October 2023

Institut d'Astrophysique de Paris
Observatoire de Paris

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